RESEARCH ARTICLE

Open Access

A hybrid twin based on machine learning enhanced reduced order model for real-time simulation of magnetic bearings



Chady Ghnatios^{1*}, Sebastian Rodriguez², Jerome Tomezyk³, Yves Dupuis³, Joel Mouterde³, Joaquim Da Silva³ and Francisco Chinesta⁴

*Correspondence:

chady.ghnatios@ensam.eu ¹PIMM Lab & ESI Chair, SKF Research Chair, Arts et Metiers Institute of Technology, Paris, France ²PIMM Lab & ESI Chair, Arts et Metiers Institute of Technology, Paris, France ³Department of Research and Development Engineering, SKF Magnetic Mechatronic, Saint-Marcel, France ⁴PIMM Lab & ESI Chair, CNRS@CREATE, Arts et Metiers Institute of Technology, Paris, France

Abstract

The simulation of magnetic bearings involves highly non-linear physics, with high dependency on the input variation. Moreover, such a simulation is time consuming and can't run, within realistic computation time for control purposes, when using classical computation methods. On the other hand, classical model reduction techniques fail to achieve the required precision within the allowed computation window. To address this complexity, this work proposes a combination of physics-based computing methods, model reduction techniques and machine learning algorithms, to tackle the requirements. The physical model used to represent the magnetic bearing is the classical Cauer Ladder Network method, while the model reduction technique is applied on the error of the physical model's solution. Later on, in the latent space a machine learning algorithm is used to predict the evolution of the correction in the latent space. The results show an improvement of the solution without scarifying the computation time. The solution is computed in almost real-time (few milliseconds), and compared to the finite element reference solution.

Keywords: Spectral method, Reduced basis, Machine learning, Magnetic bearing, Magnetic levitation, Long-short term memory

Introduction

Recent development and democratization of the available toolbox in machine learning techniques made their use increasingly popular. Thus, they become increasingly popular in engineering applications [1–4]. However, in multiple engineering systems and processes, data is expensive and hard to generate. Therefore, the ability to solely rely on data-driven techniques to generate representative models is hindered. Thus, many recent efforts aim to combine the basic principles and physical knowledge with data-driven methods [5–9]. Some of these works aim to use thermodynamics principles [10–13], while others tend to improve or correct a simulation by infusing additional information coming from data [14–16]. The result is the so called hybrid twin, or physics-aware digital twin, a combination of scientific knowledge, enhanced with machine learning techniques [17–20]. Gaussian



© The Author(s) 2024. **Open Access** This article is licensed under a Creative Commons Attribution 4.0 International License, which permits use, sharing, adaptation, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons licence, and indicate if changes were made. The images or other third party material in this article are included in the article's Creative Commons licence, unless indicated otherwise in a credit line to the material. If material is not included in the article's Creative Commons licence and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this licence, visit http://creativecommons.org/licenses/by/4.0/.

processes are also used to model data and include experimental variability into physical models [21,22].

In this work, we propose a novel digital twin building technique, based on the combination of a spectral approximation, similar to the one used in the Proper Orthogonal Decomposition method, and a time-dependent integrator based on machine learning algorithms. Similar methods already tried to combine the use of spectral decomposition with machine learning. For example, in [23], the authors built a machine learning method in the spectral domain, by aiming to change the eigenvalues and eigenvectors of the transfer operator in the direct space. This approach manipulates a matrix operator, in its eigenvalues/eigenvectors space; to adapt the solution and fit it to the measurements. In [24], the author uses a spectral decomposition of the input image before running a convolution neural network, to achieve a reduction in the memory requirements. The authors in [25,26] propose a deep learning algorithm to map or transform the reduced basis, allowing them to adapt their reduced basis to novel use cases. POD-Galerkin PINN ROM was proposed in [27] to tackle the Navier–Stokes problem. The method proposed here is different from the aforementioned ones available in the literature. The proposed method uses a weight prediction based on the basis decomposition [28]. This work however operates on the error of the quantities of interest and not the full fields. The spectral neural network takes as input the errors of the physical quantities of interest, as well as the inputs of the model, to built a correction, based on the re-weighing of the correction's reduced basis, built a priori. The correction basis is build using the POD method [29]. The aim of this work is to propose a hybrid modeling enhancing the available state of the art industrial simulation while conserving the same computation time.

The selected showcased industrial application in this work corresponds to the case of magnetic bearings. Magnetic bearings are used since several years to levitate heavy rotating shafts in critical machines. These bearings have the ability to perform quasi indefinitely since no contamination, no lubrication, no contact, and no friction are exhibited [30]. Moreover, losses by friction are totally suppressed, while rolling resistance is tremendously reduced, leading to potentially very high rotation speed [31]. Despite all these advantages, active magnetic bearings are only used on heavy equipment as they require external energy source to generate the levitation through magnetic fields. Moreover, the control of active magnetic bearings requires a feedback every few microseconds to achieve effective shaft positioning and active suppression of vibrations. The possibility of enhancing this control through simulation-based control is currently being investigated. However, such a control would require a simulation running within few microseconds. Multiple previous efforts aim to model the magnetic bearings using either analytical methods, numerical methods or data driven techniques [32-34]. Thus, this work proposes a novel technique coupling model reduction, physics based methods and advanced machine learning algorithms to tackle the simulation time requirement.

When it comes to the simulation of magnetic fields, especially in the context of rotors and stators, model reduction was used on several cases to create fast simulations [35]. However, to the best knowledge of the authors, there has been no hybrid modeling, combining scientific knowledge to data-driven techniques, in the context of creating a digital twin for active magnetic bearings simulation in real-time.

The work starts in Sect. 'Review of the used physical modeling' by reviewing the used physical modeling of the process based on the Cauer Ladder Network method [36-38]. The

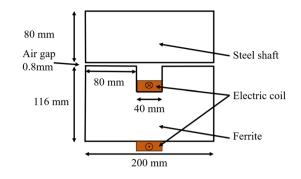


Fig. 1 A schematic drawing of the magnetic levitation bearing in use

Cauer Ladder Network method replaces a full physical-space finite element problem with a sequence of resistors and self inductance, transforming the simulation into a reduced order model problem, based on an equivalent electric circuit model to solve. Section 'Hybrid twin construction' introduces the data-driven method and the coupling between this data-driven technique and the physical model. Section 'Application to the magnetic bearings modelling's quantities of interest' shows the selected application, the results, and discusses the strong and weak points of the method. The work uses the finite element solution as the ground truth because of missing experimental measurements, the same algorithm can be used to built hybrid modeling with experimental results as ground truth when available. The work ends with some conclusions in Sect. 'Discussion and conclusions'.

Review of the used physical modeling

In this section we review the physical modeling of a magnetic bearing. The studied problem represents a single coil toy-model used to showcase the numerical technique. The schematic drawing of the coil is shown in Fig. 1. We can see in this figure the coil, which is used to create a magnetic field, leading to the levitation of the steel shaft on top of the coil. A small air gap is to be modeled with high precision to achieve a high fidelity solution.

Maxwell's equations

The problem to solve is the nonlinear magneto-quasistatic problem, a realistic assumption in the framework of magnetic bearings for industrial application, which simplifies much the full order model computation. In these settings, we consider the Maxwell's equations while neglecting the displacement current, i.e.:

$$\nabla \times \mathbf{H} = \mathbf{J}_{ind} + i\mathbf{N} \tag{1}$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \tag{2}$$

$$\nabla \cdot \mathbf{B} = 0 \tag{3}$$

$$\nabla \cdot \mathbf{J} = 0 \tag{4}$$

where **B** is the magnetic flux density, **H** the magnetic field, **J**_{ind} the eddy current density, **N** a unit current density vector and *i* is the current flowing through the coil. The constitutive laws give us that:

$$\mathbf{J}_{ind} = \sigma \mathbf{E} \tag{5}$$

$$\mathbf{B} = \mu(\mathbf{H})\mathbf{H} \tag{6}$$

Note that the electric conductivity $\sigma \neq 0$ only in the steel shaft, because we assume that the stator is laminated.

To solve this problem, we use a A^* -formulation, then a magnetic vector potential A is introduced such that:

$$\mathbf{B} = \nabla \times \mathbf{A} \text{ and } \mathbf{E} = -\frac{\partial \mathbf{A}}{\partial t}.$$
 (7)

In order to impose a voltage in the electrical coil winding, we consider the circuit equation:

$$\frac{d\phi}{dt} + Ri = U \text{ with } \phi = \int_{\Omega_{coil}} \mathbf{A} \cdot \mathbf{N} \, dx, \tag{8}$$

where *R* is the resistance of the coil and ϕ the magnetic flux. The complete problem is then written as:

$$\begin{cases} \nabla \times \left(\frac{1}{\mu} \nabla \times \mathbf{A}\right) + \sigma \frac{\partial \mathbf{A}}{\partial t} - i\mathbf{N} = 0\\ \frac{d\phi}{dt} + Ri = U \end{cases}$$
(9)

U being the imposed electric tension of the magnetic bearing's coil. As we are addressing a 2D in-plane problem for the sake of simplicity, the potential **A** is reduced to a scalar representing only the out of plane component of **A**. At the boundary of the domain, we impose a homogeneous Dirichlet boundary condition.

The electromagnetic force is computed using the Maxwell's tensor Eq. [39], i.e.:

$$Force = \int_{\Gamma_{Force}} \frac{1}{2\mu_0} \|\mathbf{B}\|^2 \mathbf{n} + \frac{1}{\mu_0} (\mathbf{B} \cdot \mathbf{n}) \mathbf{B} \, dx,\tag{10}$$

where the line domain Γ_{Force} (included in the air) encloses the moving part, **n** is the normal to this line. Finally, μ_0 is the magnetic permeability of the empty space.

Constructing a reference solution using finite element

In the first part of this work, we use the finite element method (FE) to solve the Eq. (9), which will lead us to a reference solution. The reference solution is used later on as our target solution. We can see in Fig. 2 the adaptive mesh used for the FE discretization, which is refined near the airgap and in the conductive part to take into account the eddy currents.

The FE solution is computed in the 2D domain in a square of dimensions 700×700 *mm*, with an air box of 250 *mm* around the simulated magnetic bearing.

The transient finite element solution is computed for 14 different time variations of the input electric potential U, using 14 different control use cases. We know that PWM exacerbates the eddy current, then to have some example with many eddy currents we implement a simple p-controller with saturation to control the voltage and get as a set point the magnetic flux.

The available use-cases are illustrated in table 1. One of the finite elements solutions is shown in Fig. 3, for illustrative purposes. In 3(a) we illustrate the out of plane component of the magnetic vector potential \mathbf{A}_z , which, for sake of simplicity, will be referred as *A* for the rest of this work. Moreover, in 3(b) we illustrate the magnitude of the magnetic flux density $\|\mathbf{B}\|$.

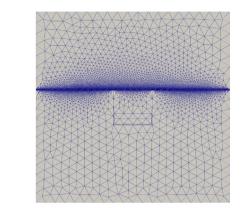
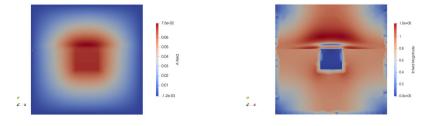


Fig. 2 The adopted mesh considered for the high fidelity finite element simulation of the magnetic levitation bearing

z ¥

Use case	Tension Amplitude	Shape form	Frequency	Train/Test
1	40V	Ramp	None	Train
2	40V	Sin	1Hz	Train
3	40V	Sin	5Hz	Train
4	40V	Step	None	Train
5	70V	Vdc	3Hz	Train
6	100V	Double step	None	Train
7	100V	Vdc	6Hz	Train
8	100V	Vdc+step	None	Train
9	300V	Double step	None	Train
10	300V	Vdc	3Hz	Train
11	100V	Vdc+slope	3Hz	Test
12	300V	Vdc	5Hz	Test
13	300V	Vdc+slope	6Hz	Test
14	300V	Vdc+step	None	Test



(a) A sample finite element solution for A (b) A sample finite element solution for $\|\mathbf{B}\|$

Fig. 3 A finite element solution for a selected use case solved in a time domain $t \in [0; 5]s$, and illustrated at t = 5s in this figure

A fast physics-based modeled solution

Recent developments in the simulation of magnetic fields generated by coils have lead to the creation of a novel simulation technique based on the use of spectral, mode shape based, representation. This method is currently known by the name of Cauer Ladder Network Method [36-38]. The method consists in assuming that *A* can be written in a separated form using:

$$A = \sum_{q=1}^{n_q} c_q \mathbf{a}_q(x, y) \tag{11}$$

where n_q is the number of selected layers in the ladder Network representation, c_q is a constant to compute, equal to the current in the equivalent circuit, and \mathbf{a}_q are the orthogonal basis representation vectors. In this work, the presented results are shown for $n_q = 8$, as beyond $n_q = 8$ it is observed that the results do not improve anymore in the selected use case. To compute \mathbf{a}_q and c_q we use the matrix formulation method presented in [38]. The method produces an electric circuit of several (*R*, *L*) branches, where *R* is a resistance and *L* denotes an inductance. The current intensity in each branch of the ladder Network is noted c_j and will be used later as input for the correction of the results. An example of the equivalent Cauer ladder network is shown in Fig. 4.

We note:

$$K = \int v(\mathbf{A}_{\text{REF}}) \nabla \times w_i \cdot \nabla \times w_j \, dx$$

$$S = \int \sigma w_i w_j \, dx \tag{12}$$

where A_{REF} corresponds to the solution of a magnetostatic problem with i = 15 (that corresponds to the first equation of (9) with $\sigma = 0$). In fact, this term is used to construct the Cauer Ladder Network, that is linear, near a saturation operating point, in order to obtain a linear model as close as possible to the exact nonlinear model. Algorithm 1 describes how he Cauer network is built.

Algorithm 1: Algorithm for determining the coefficients of the Cauer network

Result: a_i, L_i, R_i for $1 \le i \le n_q$ $R_1 = 1$ $a_1 = K^{-1}N$ $L_1 = a_1^t K a_1$ $e_1 = 0$ **for** $i = 2: n_q$ **do** $e_{i+1} = e_i - \frac{a_i}{L_i}$ $R_{i+1} = \frac{1}{e_{i+1}^t S e_{i+1}}$ $a_{i+1} = K^{-1} (R_{i+1} S e_{i+1} + K a_i)$ $L_{i+1} = a_{i+1}^t K a_{i+1}$ **end** R_1 = resistor of the coil

The Cauer Ladder Network solution of the selected use case is shown in Fig. 5 when using eight ladder circuits (i.e. $n_q = 8$). Is it noted that further increase of the number of Cauer Ladder circuits n_q would no longer improve the solution considerably. We can clearly note a similar trend as for the finite element solution shown in Fig. 3, especially for the *A* field, but the *B* field seems to show large differences, especially when it comes to the in-depth penetration of the magnetic field.

The Cauer Ladder Network Method is an excellent method to simulate linear problems and linear dependencies. However, in highly nonlinear problems exhibiting large eddy currents and magnetic saturation, like the one showcased in this work, its prediction power degrades dramatically with respect to the reference solution.

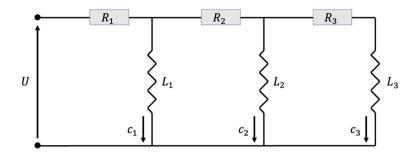
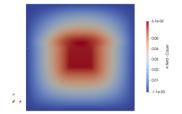
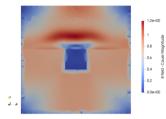


Fig. 4 A equivalent Cauer ladder network with $n_a = 3$



(a) A sample Cauer Ladder Network Method solution for ${\cal A}$



(b) A sample Cauer Ladder Network Method solution for $\|\mathbf{B}\|$

Fig. 5 A Cauer Ladder Network Method solution for the same selected use case shown in Fig. 3, computed for $t \in [0; 5]s$ and illustrated at t = 5s

Hybrid twin construction

This section aims at building-up a hybrid twin capable of improving the fast physics-based model presented in Sect. 'A fast physics-based modeled solution'. The main driving reason to consider an improvement of the Cauer Ladder Network Method is its ability to capture a large part of the solution within a fraction of a second per time step, using a standard PC.

To start the hybrid modeling of the magnetic levitation, we first compute a reduced basis of the error, using only 10 out of the 14 available simulations in the database. To compute the reduced basis, we use the traditional Proper Orthogonal Decomposition technique. Thus, the solution errors in the physical space (x, y), are placed, at every time-step, as a vectors in a large matrix of error snapshots **X**. For instance:

$$\mathbf{C}_{j} = A_{j}^{FEM} - A_{j}^{Cauer}, j = 1, \cdots, 6$$

$$\mathbf{X} = \begin{bmatrix} \vdots & \vdots & \vdots \\ C_{1} & C_{2} & \cdots & C_{6} \\ \vdots & \vdots & \vdots \end{bmatrix}.$$
(13)

with A_j^{FEM} the finite element method's solution at time step *j* and A_j^{Cauer} the Cauer Ladder Network's one. Later on, a singular value decomposition (*SVD*) is used to select the most relevant singular vectors of the solution. For instance, we can note:

$$[\mathbf{V}, \mathbf{\Lambda}, \mathbf{T}] = SVD(\mathbf{X}) \tag{14}$$

V is the matrix of space singular vectors, the diagonal of the matrix Λ contains the sorted decreasing singular values λ_k , while **T** contains the time vector basis. The selected reduced basis vectors V_k of space variation of the solution are selected using all k satisfying:

$$\frac{\lambda_k}{\lambda_1} > \epsilon,$$
 (15)

where ϵ is a tolerance value. In this use case, the tolerance is set to $\epsilon = 10^{-4}$. This tolerance has lead to the use of approximately 120 vectors V_i in the correction's reduced space. The correction term **C** is now supposed to have the form:

$$\mathbf{C} = \sum_{k=1}^{n} \alpha_k(t) V_k,\tag{16}$$

with n = 120 the number of vectors in the reduced basis, and α_i are the reduced coordinates, or the weights of the reduced basis vectors V_k .

To efficiently compute α_i within a very limited time, and possibly in a non-incremental manner, α_i are supposed to be the output of a surrogate model \mathcal{H} built using neural networks. To take into consideration the memory effect of the solution, depending on previous inputs/outputs of the system, the surrogate model \mathcal{H} uses a memory layer. For instance, the time dependency of the solution on previous inputs/outputs combination is accounted for using a Long-short Term Memory layer (LSTM layer) [40]. The use of LSTM requires the input to have several time-steps, and not only the current one. The inputs of the surrogate model are in fact the imposed electrical tension U, the computed Cauer Ladder Network currents in the eight branches c_j for $j = 1, \dots, 8$, the Cauer computed force F, the magnetic flux ϕ , and the time t. The force F and the magnetic flux ϕ are our global quantities of interest computed from a post-processing of the Cauer Ladder Network's solution A^{Cauer} . All these inputs are given in the form of the last κ time-steps, such as at every time step i we can write:

$$\boldsymbol{\alpha}_{i} = \mathcal{H}\left(\begin{pmatrix} U_{i} \\ U_{i-1} \\ \vdots \\ U_{i-\kappa+1} \end{pmatrix}, \begin{pmatrix} F_{i} \\ F_{i-1} \\ \vdots \\ F_{i-\kappa+1} \end{pmatrix}, \begin{pmatrix} \phi_{i} \\ \phi_{i-1} \\ \vdots \\ \phi_{i-\kappa+1} \end{pmatrix}, \begin{pmatrix} t_{i} \\ t_{i-1} \\ \vdots \\ t_{i-\kappa+1} \end{pmatrix}, \begin{pmatrix} c_{i}^{8} \\ c_{i-1}^{8} \\ \vdots \\ c_{i-\kappa+1}^{8} \end{pmatrix} \right)$$

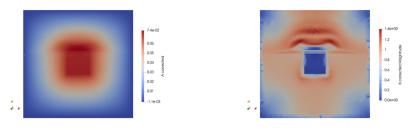
$$(17)$$

Using this notation, we accept that the first κ time steps will not be corrected, and the correction will only start after κ time steps are computed. In the selected surrogate model \mathcal{H} , a single dense fully connected layer follows the LSTM layer and outputs the *n* required values in the α vector, at every time step. One should note that the inputs of the surrogate model depend only on the Cauer Ladder Network solution and not on the correction itself. Thus, all α_i for all time steps can be computed simultaneously.

The surrogate model \mathcal{H} is named a spectral network. In fact, as previously explained, it uses the spectral decomposition of the field, through a singular value decomposition, and reconstructs the solution in a form of a reduced order model. The weights are predicted as a function of the current and previous time inputs, while fixing the reduced basis of the approximation **V**.

Remark 1 The corrected solution is finally computed such as:

$$A^{corrected} = A^{Cauer} + \sum_{k=1}^{n} \alpha_k(t) V_k, \tag{18}$$



(a) A sample Cauer Ladder Network Method corrected using the spectral network approach. Figure showing A fields.

(b) A sample Cauer Ladder Network Method corrected using the spectral network approach. Figure showing $\|\mathbf{B}\|$

Fig. 6 Corrected hybrid model solution for the same selected use case shown in Fig. 3, illustrated at t = 5s

however, since $\alpha_k(t)$ is nonlinear as a function of global quantities computed using A^{Cauer} , the correction is nonlinear as a function of the Cauer Ladder Network's solution.

To train the network, the weights and biases of the two selected layers can be computed using a gradient descent algorithm and the following loss function \mathcal{L} :

$$\mathcal{L} = \sum_{j=1}^{6} \left(\sum_{k=1}^{n} \alpha_{k}^{j}(t) V_{k} - \mathbf{C}_{j} \right)^{2} + \beta \sum_{w=1}^{m} W_{w}^{2},$$
(19)

where W are the weights to be optimized in the neural network and β a regularization coefficient. The optimization is performed using ADAM optimization algorithm with a custom adaptive learning rate. The gradients are computed using automatic differentiation in *Keras*–*Tensorflow* toolbox, using tensorflow's GradientTape functionality. The training is only performed on the selected 10 data sets, out of the 14 available simulations.

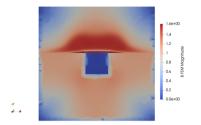
The solution of the corrected fields $A^{corrected}$ for $\kappa = 5$ is shown in Fig. 6 for the same use case illustrated in Fig. 3. The solution shows a good improvement with respect to the Cauer Ladder Network solution. One should note that the solution shown in Fig. 6 belongs to a training set. Another solution belonging to the testing set is illustrated in Fig. 7, which shows simultaneously the finite element solution in 7a, the Cauer Ladder Network solution in Fig. 7b, and the corrected solution in Fig. 7c. The testing solution shows the amplitude of the magnetic field $\|\mathbf{B}\|$, and also showcases a large improvement of the solution after correction. This improvement is more tangible when exploring global quantities of interest, as illustrated in sect. 'Application to the magnetic bearings modeling's quantities of interest'.

One should note that the computation of the solution's weights and correction, for all time steps, was performed within 1 to 1.1 s, for any of the 14 available use cases.

The relative error maps on the testing set are shown in Fig. 8. The error for the Cauer Ladder Network reaches more than 50% on the *A* fields. The error of the corrected Cauer by the proposed Spectral Net does not exceed 11%. Moreover, the error is localized away from the region of interest, which is mainly around the coil and shaft separation region.

Application to the magnetic bearings modeling's quantities of interest

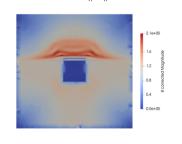
In this section we illustrate the global quantities for all the training and testing sets, as a function of time. The global quantities of interest are the force *F* applied on the steel shaft and the magnetic flux ϕ .



- 40+00 - 35 - 2 - 2 - 36 - 15 - 10 - 0.6 - 0.09+00

(b) Cauer Ladder Network's solution showing

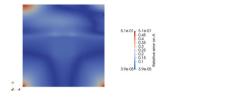
(a) Finite element solution showing $\|\mathbf{B}\|$



 $\|\mathbf{B}\|$

(c) Hybrid twin corrected solution showing $\|\mathbf{B}\|$

Fig. 7 A solution from the testing set at the same time step. We remind the reader to mind the axis



, 2 ,

(a) Relative error on the A fields computed by the Cauer Ladder Network



(b) Relative error on the A fields computed by the corrected Cauer Ladder Network using the Spectral Net method

Fig. 8 Relative errors of a solution from the testing set at the same time step

Figures 9 and 10 illustrate the global quantities for the training sets, for the forces and the magnetic flux respectively. Figs. 9 and 10 also compare the finite element solution, the Cauer Ladder Network's solution, and the corrected hybrid twin solution. in all the figures E_m stands for the mean relative error in percentage point-wise, as defined in Eq. (20), and E_{max} is the maximum relative error in percentage, as defined in Eq. (21):

$$E_m = \frac{100}{K} \sum_{i=1}^{K} \frac{\|\hat{Z}_i - Z_i\|_2}{\|Z_i\|_2}$$
(20)

$$E_{max} = 100 \times \max\left\{\frac{\|\hat{Z}_i - Z_i\|_2}{\|Z_i\|_2}\right\}_{i=1,\cdots,K}$$
(21)

Z being the quantity of interest and K the number of time steps available in the sequence. The testing set solutions, post processed to compute the global quantities of interest, are illustrated in Figs. 11 and 12 for the force and the flux respectively. The testing set shows an excellent agreement of the hybrid twin solution with the finite element one, and a tremendous improvement when compared to the Cauer Ladder Network results.

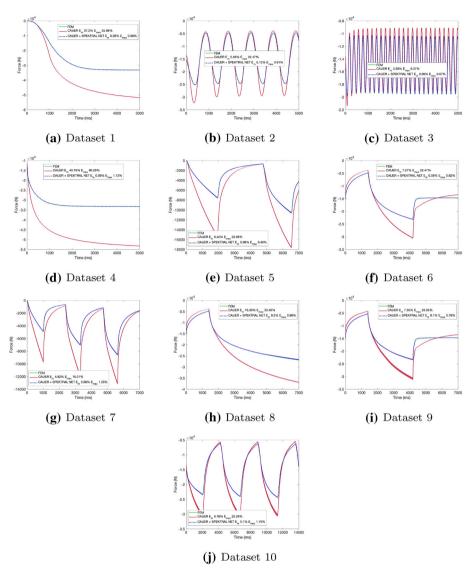


Fig. 9 Post processing of the force F in the training set

An effort to model the global quantities was taken by the authors using fully data-driven techniques, using many advanced machine learning techniques. When not infusing physical information in the model, the solution remains unstable and couldn't be generalized to different scenarios or use cases. The authors tested an LSTM method, a Resnet with embedded LSTM to compute the derivatives, as well as a NeuralODE, based on fully data-driven approximation of the derivatives. However, none of the tested techniques was able to produce stable results for thousands of time steps, within a fraction of a second. The presented results prove the robustness of the proposed approach. Many of the testing sets were performed at a different maximum tension *U*, different scenario schemes, and different frequencies, some never seen in the training zone. The algorithm can perform in extrapolation with very reasonable outputs, while computing the results of about 14,000 time steps within about 1 s on a usual PC.

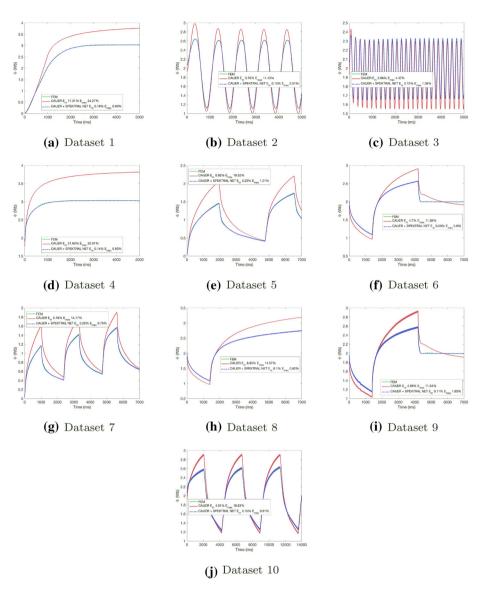


Fig. 10 Post processing of the magnetic flux ϕ in the training set

The results show excellent accuracy while the simulation performs in almost real-time. When it comes to control, the inference of a single time step is currently about 200 ms using Python libraries, limited by the Cauer ladder Network computation requirements. Nevertheless, the simulation of few thousand time steps takes about 1 to 1.1 s, leveraging the vector and tensor computation possibilities. The only calculation performed online is the prediction of the basis weights α_i using the neural network at every time step *i*.

The required computation time in the electronic control is less a feedback frequency of 2000 Hz, that is around 0.5*ms* per feedback. The resultant algorithm is not implemented on the electronic material yet. A speedup of 100 times between python computing and PCB computing is realistic and achieved in multiple previous applications. The algorithm achieves the required computation time on PC. The code optimization and the use of dedicated PCB material achieves the required computation time for control purposes.

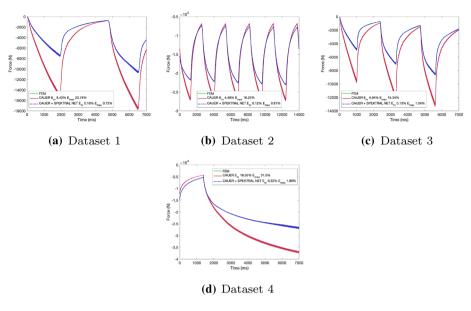


Fig. 11 Post processing of the force F in the testing set

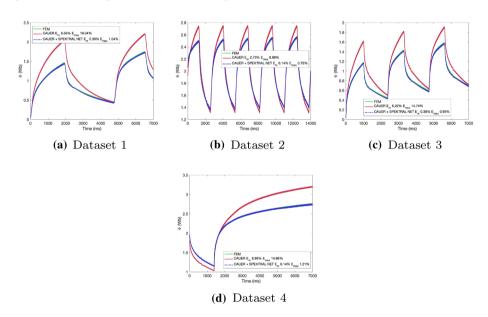


Fig. 12 Post processing of the magnetic flux ϕ in the testing set

Discussion and conclusions

This work proposes a hybrid modeling, enhancing the available state of the art simulation of magnetic field propagation, without impacting computation time. The results show a relevant enhancement of the solution throughout all the simulated time domain. Moreover, the solution is stable and reliable, as depicted in the different testing cases. The Cauer + Specktral Net solution enhanced the Cauer Ladder Network solution, while preserving stability and computation time.

In this work, we illustrated a novel machine learning hybrid twin architecture, which relies first on the use of spectral decomposition of a solution, in a form of a reduced order basis. Later on, the method finds the solution weights using a machine learning algorithm.

The proposed spectral network improves a physics-based simulation by computing a correction, using a reduced basis of the error with respect to a ground truth, high fidelity, finite element solution. This work uses the finite element solution as the ground truth values, since experimental measurements are missing. The same algorithm can be used to build hybrid modeling corrections fitting the experimental measurements when available.

Moreover, the algorithm uses explainable physics, and its limitations can be understood, and controlled to a certain extent, using the basic principles of electromagnetic.

Acknowledgements

This work was supported by the SKF research chair at Arts et Métiers institute of technology-Paris.

Author contributions

All the authors participated equally in the paper content.

Availability of data and materials

Interested reader can contact authors.

Declarations

Competing interests

The authors declare that they have no competing interests.

Received: 21 August 2023 Accepted: 18 January 2024 Published online: 03 February 2024

References

- 1. Wright LG, Onodera T, Stein MM, Wang T, Schachter DT, Hu Z, McMahon PL. Deep physical neural networks trained with backpropagation. Nature. 2022;601(7894):549–55. https://doi.org/10.1038/s41586-021-04223-6.
- Daniel T, Casenave F, Akkari N, Ryckelynck D. Model order reduction assisted by deep neural networks (rom-net). Adv Model Simul Eng Sci. 2020;7(1):16. https://doi.org/10.1186/s40323-020-00153-6.
- 3. Champaney V, Amores VJ, Garois S, Irastorza-Valera L, Ghnatios C, Montàns FJ, Cueto E, Chinesta F. Modeling systems from partial observations. Front Mater. 2022. https://doi.org/10.3389/fmats.2022.970970.
- Ghnatios C, Champaney V, Pasquale A, Chinesta F. A regularized real-time integrator for data-driven control of heating channels. Computation. 2022. https://doi.org/10.3390/computation10100176.
- Soize C, Farhat C. A nonparametric probabilistic approach for quantifying uncertainties in low- and high-dimensional nonlinear models. Int J Numer Meth Eng. 2016;109:837–88. https://doi.org/10.1002/nme.5312.
- Farhat C, Bos A, Avery P, Soize C. Modeling and quantification of model-form uncertainties in eigenvalue computations using a stochastic reduced model. Am Inst Aeronaut Astronaut J. 2017;56(3):1–22.
- Soize C, Farhat C. Probabilistic learning for modeling and quantifying model-form uncertainties in nonlinear computational mechanics. Int J Numer Meth Eng. 2019;117:819–43. https://doi.org/10.1002/nme.5980.
- Azzi M-J, Ghnatios C, Avery P, Farhat C. Acceleration of a physics-based machine learning approach for modeling and quantifying model-form uncertainties and performing model updating. J Comput Inf Sci Eng. 2022;23(1): 011009. https://doi.org/10.1115/1.4055546.
- Ghnatios C, Barasinski A, Cueto E. Editorial: advanced materials modeling combining model order reduction and data science. Front Mater. 2023;9:1096233. https://doi.org/10.3389/fmats.2022.1096233.
- 10. Moya B, Badías A, González D, Chinesta F, Cueto E. A thermodynamics-informed active learning approach to perception and reasoning about fluids. Comput Mech. 2023;72(3):577–91. https://doi.org/10.1007/s00466-023-02279-x.
- 11. Pichi F, Moya B, Hesthaven JS. A graph convolutional autoencoder approach to model order reduction for parametrized PDEs. 2023. 2305.08573.
- Gonzàlez D, Chinesta F, Cueto E. Thermodynamically consistent data-driven computational mechanics. Contin Mech Thermodyn. 2019;31:239–53.
- Ghnatios C, Alfaro I, Gonzàlez D, Chinesta F, Cueto E. Data-driven generic modeling of poroviscoelastic materials. Entropy. 2019. https://doi.org/10.3390/e21121165.
- Li H, Yang D, Cao H, Ge W, Chen E, Wen X, Li C. Data-driven hybrid petri-net based energy consumption behaviour modelling for digital twin of energy-efficient manufacturing system. Energy. 2022;239: 122178. https://doi.org/10. 1016/j.energy.2021.122178.
- 15. Wu J, Ghnatios C, Mordillat P, Tourbier Y, Chinesta F. Functional parametric elasto-dynamics for efficient multicomponent design. Computation. 2022. https://doi.org/10.3390/computation10120218.
- Ghnatios C, Gérard P, Barasinski A. An advanced resin reaction modeling using data-driven and digital twin techniques. Int J Mater Form. 2022;16(1):5. https://doi.org/10.1007/s12289-022-01725-0.
- Wang Y, Zhong L. Nas-pinn: Neural architecture search-guided physics-informed neural network for solving pdes. J Comput Phys. 2024;496: 112603. https://doi.org/10.1016/j.jcp.2023.112603.
- Saqlain S, Zhu W, Charalampidis EG, Kevrekidis PG. Discovering governing equations in discrete systems using pinns. Commun Nonlinear Sci Numer Simul. 2023;126: 107498. https://doi.org/10.1016/j.cnsns.2023.107498.
- Raissi M, Perdikaris P, Karniadakis GE. Physics-informed neural networks: a deep learning framework for solving forward and inverse problems involving nonlinear partial differential equations. J Comput Phys. 2019;378:686–707. https:// doi.org/10.1016/j.jcp.2018.10.045.

- Ghnatios C, di Lorenzo D, Champaney V, Cueto E, Chinesta F. Optimal velocity planning based on the solution of the Euler-Lagrange equations with a neural network based velocity regression. Discrete Continuous Dyn Syst. 2023. https://doi.org/10.3934/dcdss.2023080.
- Wahlström N, Kok M, Schön TB, Gustafsson F. Modeling magnetic fields using Gaussian processes. In: 2013 IEEE international conference on acoustics, speech and signal processing, pp. 3522–3526. 2013. https://doi.org/10.1109/ ICASSP.2013.6638313
- 22. Sun Y, Wang S, Lu Y, Xu J. Gaussian process dynamic modeling and backstepping sliding mode control for magnetic levitation system of maglev train. J Theor Appl Mech. 2022; 60(1): 49–62. https://doi.org/10.15632/jtam-pl/143676
- Giambagli L, Buffoni L, Carletti T, Nocentini W, Fanelli D. Machine learning in spectral domain. Nat Commun. 2021;12(1):1330. https://doi.org/10.1038/s41467-021-21481-0.
- Guan B, Zhang J, Sethares WA, Kijowski R, Liu F. SpecNet: spectral domain convolutional neural network. arXiv. 2019. https://doi.org/10.48550/ARXIV.1905.10915.
- Fresca S, Manzoni A. Pod-dl-rom: enhancing deep learning-based reduced order models for nonlinear parametrized pdes by proper orthogonal decomposition. Comput Methods Appl Mech Eng. 2022;388: 114181. https://doi.org/10. 1016/j.cma.2021.114181.
- Papapicco D, Demo N, Girfoglio M, Stabile G, Rozza G. The neural network shifted-proper orthogonal decomposition: a machine learning approach for non-linear reduction of hyperbolic equations. Comput Methods Appl Mech Eng. 2022;392: 114687. https://doi.org/10.1016/j.cma.2022.114687.
- Hijazi S, Freitag M, Landwehr N. Pod-Galerkin reduced order models and physics-informed neural networks for solving inverse problems for the navier-stokes equations. Adv Model Simul Eng Sci. 2023;10(1):5. https://doi.org/10.1186/ s40323-023-00242-2.
- Hesthaven JS, Ubbiali S. Non-intrusive reduced order modeling of nonlinear problems using neural networks. J Comput Phys. 2018;363:55–78. https://doi.org/10.1016/j.jcp.2018.02.037.
- Carlberg K, Farhat C. A low-cost, goal-oriented 'compact proper orthogonal decomposition' basis for model reduction of static systems. Int J Numer Meth Eng. 2011;86(3):381–402.
- Wajnert D, Tomczuk B. Two models for time-domain simulation of hybrid magnetic bearing' characteristics. Sensors. 2022. https://doi.org/10.3390/s22041567.
- Xu C, Zhao F, Xu G, Lian G. General-purpose controller for multiple types of magnetic bearing rotor systems. Energy Rep. 2023;9:345–54. https://doi.org/10.1016/j.egyr.2022.10.423.
- Sikora BM, Piłat AK. Analytical modeling and experimental validation of the six pole axial active magnetic bearing. Appl Math Model. 2022;104:50–66. https://doi.org/10.1016/j.apm.2021.10.024.
- 33. Taha OW, Hu Y. Modeling of a digital twin for magnetic bearings. Appl Sci. 2023. https://doi.org/10.3390/app13148534.
- Yoo SJ, Kim S, Cho K-H, Ahn H-J. Data-driven self-sensing technique for active magnetic bearing. Int J Precis Eng Manuf. 2021;22(6):1031–8. https://doi.org/10.1007/s12541-021-00525-x.
- Sancarlos A, Ghnatios C, Duval J-L, Zerbib N, Cueto E, Chinesta F. Fast computation of multi-parametric electromagnetic fields in synchronous machines by using pgd-based fully separated representations. Energies. 2021. https://doi. org/10.3390/en14051454.
- Eskandari H, Matsuo T. Comparison study of first-order approximations of nonlinear Eddy-current field using Cauer ladder network method. IEEE Trans Magn. 2021;57(6):1–4. https://doi.org/10.1109/TMAG.2021.3060503.
- 37. Shindo Y, Kameari A, Sugahara K, Matsuo T. Dynamical model of an electromagnet using Cauer ladder network representation of Eddy-current fields. IEEJ J Ind Appl. 2018;7(4):305–13. https://doi.org/10.1541/ieejjia.7.305.
- Matsuo T, Kameari A, Sugahara K, Shindo Y. Matrix formulation of Cauer ladder network method for efficient eddycurrent analysis. In: 2018 IEEE International Magnetics Conference (INTERMAG), pp. 1–1 (2018). https://doi.org/10. 1109/INTMAG.2018.8508788
- Bastos JPA, Sadowski N. Electromagnetic modeling by finite element methods, vol. 1. Boca Raton, FL: CRC Press; 2003. https://doi.org/10.1201/9780203911174.
- Hochreiter S, Schmidhuber J. Long short-term memory. Neural Comput. 1997;9(8):1735–80. https://doi.org/10.1162/ neco.1997.9.8.1735.

Publisher's Note

Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.